

INFORMATION ON MASTER'THESIS

1. Full name: Nguyen Duc Vinh
2. Sex: Male
3. Date of birth: 07.02.1978
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5. Decision to recognize students of :2714QĐ-CTSV on May 18/12/2008
6. The changes in the training process: No
7. Title of dissertation:
"To the pair Squark Susy-QCD in e^+e^- processes with complex parameters,"
8. Major: Theoretical Physics and Mathematical Physics 9. Code: 60.44.01
10. Scientific staff guide:
TS. Pham Thuc Tuyen - Physical Sciences - University of Natural Sciences
11. Summary results of the thesis:

Interaction Lagrangian and Feynman rules in the MSSM

To obtain the mass spectrum of particles in a physical theory we have carried out standard procedures symmetry breaking with the average value of the Higgs vacuum. I will choose the vacuum average of the two Higgs multi-line as follows:

to satisfy the equation : $\langle H^1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} v_1 \\ 0 \end{pmatrix}$, $\langle H^2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix}$

v_1, v_2 to satisfy the equation

$$\begin{cases} \left[\frac{e^2}{8 \sin^2 \theta \cos^2 \theta} (v_1^2 - v_2^2) + m_{H_1}^2 + |\mu|^2 \right] v_1 = -\mu_s v_2 \\ \left[-\frac{e^2}{8 \sin^2 \theta \cos^2 \theta} (v_1^2 - v_2^2) + m_{H_2}^2 + |\mu|^2 \right] v_2 = -\mu_s v_1 \end{cases}$$

Charge e of the particles related to the coefficients related $g_{1,2}$ through a parameter θ called the Weinberg angle. $e = g_1 \cos \theta = g_2 \sin \theta$

Lagrangian interaction between the quark and the gauge field photon (γ), wion, Zion and gluon (g):

$$\begin{aligned}
L_{qq\gamma} &= -ee_q \bar{q} \gamma^\mu q A_\mu \\
L_{qqZ} &= -\frac{g}{\cos\theta_w} \bar{q} \gamma^\mu \left\{ (I_{qL}^3 - e_q \sin^2 \theta_w) P_L - e_q \sin^2 \theta_w P_R \right\} q Z_\mu \\
&= -\frac{g}{\cos\theta_w} \bar{q} \gamma^\mu (C_{qL} P_L + C_{qR} P_R) q \\
C_{qL,R} &= I_{qL,R}^3 - e_q \sin^2 \theta_w \\
L_{qqW} &= -\frac{g}{\sqrt{2}} (W_\mu^+ \bar{t} \gamma^\mu P_L b + W_\mu^- \bar{b} \gamma^\mu P_L t) \\
L_{qqg} &= -g_s T_{rs}^a G_\mu^a \bar{q}_r \gamma^\mu q_s
\end{aligned}$$

Interaction between the super-partner of the quark (scalar quark) with the standard. They are:

- Squark-squark-photon

$$\begin{aligned}
L_{\tilde{q}\tilde{q}\gamma} &= iee_q \left(\tilde{q}_L^* \tilde{\partial}^\mu \tilde{q}_L + \tilde{q}_R^* \tilde{\partial}^\mu \tilde{q}_R \right) A_\mu \\
&= iee_q A_\mu \left(R_{i1}^{\tilde{q}} R_{j1}^{\tilde{q}} + R_{i2}^{\tilde{q}} R_{j2}^{\tilde{q}} \right) \tilde{q}_j^* \tilde{\partial}^\mu \tilde{q}_i \\
&= iee_q \delta_{ij} A_\mu \tilde{q}_j^* \tilde{\partial}^\mu \tilde{q}_i
\end{aligned}$$

Squark-squark- Z^0

$$L_{\tilde{q}\tilde{q}Z} = \frac{ig}{\cos\theta_w} Z_\mu \left(C_{qL} \tilde{q}_L^* \tilde{\partial}^\mu \tilde{q}_L + C_{qR} \tilde{q}_R^* \tilde{\partial}^\mu \tilde{q}_R \right) = \frac{ig}{\cos\theta_w} c_{ij} Z_\mu \tilde{q}_j^* \tilde{\partial}^\mu \tilde{q}_i$$

- Squark-squark- W^\pm

$$L_{\tilde{q}\tilde{q}W} = \frac{ig}{\sqrt{2}} \left(W_\mu^+ \tilde{t}_L^* \tilde{\partial}^\mu \tilde{b}_L + W_\mu^- \tilde{b}_R^* \tilde{\partial}^\mu \tilde{t}_R \right) = \frac{ig}{\sqrt{2}} \left(R_{i1}^{\tilde{q}} R_{j1}^{\tilde{q}} W_\mu^+ \tilde{t}_j^* \tilde{\partial}^\mu \tilde{b}_i + R_{i2}^{\tilde{q}} R_{j2}^{\tilde{q}} W_\mu^- \tilde{b}_j^* \tilde{\partial}^\mu \tilde{t}_i \right)$$

- Squark-squark-gluon

$$L_{\tilde{q}\tilde{q}g} = ig_s T_{rs}^a G_\mu^a \left(\tilde{q}_{Lr}^* \tilde{\partial}^\mu \tilde{q}_{Ls} + \tilde{q}_{Rr}^* \tilde{\partial}^\mu \tilde{q}_{Rs} \right) = ig_s T_{rs}^a \delta_{ij} G_\mu^a \tilde{q}_{jr}^* \tilde{\partial}^\mu \tilde{q}_{is}$$

The interaction between the squark with the five Higgs field in:

$$L_{qqH} = s_1^q h^0 \bar{q} q + s_2^q H^0 \bar{q} q + s_3^q A^0 \bar{q} \gamma^5 q + H^+ \bar{t} \left(s_4^a P_L + s_4^b P_R \right) b + H^- \bar{b} \left(s_4^b P_L + s_4^a P_R \right) t$$

Interaction between the squark and Higgs boson can be written in general as

$$\text{follows: } L_{\tilde{q}\tilde{q}H} = H_k \left(\tilde{q}_L^{\beta*}, \tilde{q}_R^{\beta*} \right) \hat{G}_k^\alpha \begin{pmatrix} \tilde{q}_L^\alpha \\ \tilde{q}_R^\alpha \end{pmatrix} = \left(G_k^\alpha \right)_{ij} H_k \tilde{q}_j^{\beta*} \tilde{q}_i^\alpha$$

Quark-squark-chargino

$$\begin{aligned}
L_{q\tilde{q}\tilde{\chi}^+} &= g\bar{t}(-U_{1j}P_R + Y_t V_{2j}P_L)\tilde{\chi}_j^+\tilde{b}_L + g\bar{t}(Y_b V_{2j}P_R)\tilde{\chi}_j^+\tilde{b}_R \\
&+ g\bar{b}(-V_{1j}P_R + Y_b U_{2j}P_L)\tilde{\chi}_j^{+c}\tilde{t}_L + g\bar{b}(Y_t V_{2j}P_R)\tilde{\chi}_j^{+c}\tilde{t}_R \\
&+ g\tilde{\chi}_j^+(-U_{1j}P_L + Y_t V_{2j}P_R)t\tilde{b}_L^* + g\tilde{\chi}_j^+(Y_b V_{2j}P_L)t\tilde{b}_R^* \\
&+ g\tilde{\chi}_j^{+c}(-V_{1j}P_L + Y_b U_{2j}P_R)b\tilde{t}_L^* + g\tilde{\chi}_j^{+c}(Y_t V_{2j}P_L)b\tilde{t}_R^* \\
&= g\bar{t}(l_{ij}^b P_R + k_{ij} P_L)\tilde{\chi}_j^+\tilde{b}_L + g\bar{t}(Y_b V_{2j}P_R)\tilde{\chi}_j^+\tilde{b}_R \\
&+ g\bar{b}(-V_{1j}P_R + Y_b U_{2j}P_L)\tilde{\chi}_j^{+c}\tilde{t}_L + g\bar{b}(Y_t V_{2j}P_R)\tilde{\chi}_j^{+c}\tilde{t}_R \\
&+ g\tilde{\chi}_j^+(-U_{1j}P_L + Y_t V_{2j}P_R)t\tilde{b}_L^* + g\tilde{\chi}_j^+(Y_b V_{2j}P_L)t\tilde{b}_R^* \\
&+ g\tilde{\chi}_j^{+c}(-V_{1j}P_L + Y_b U_{2j}P_R)b\tilde{t}_L^* + g\tilde{\chi}_j^{+c}(Y_t V_{2j}P_L)b\tilde{t}_R^*
\end{aligned}$$

Quark-squark-neutralino

$$\begin{aligned}
L_{q\tilde{q}\tilde{\chi}^0} &= g\bar{q}(f_{Lk}^q P_R + h_{Lk}^q P_L)\tilde{\chi}_k^0\tilde{q}_L + g\bar{q}(h_{Rk}^q P_R + f_{Rk}^q P_L)\tilde{\chi}_k^0\tilde{q}_R + H.c \\
&= g\bar{q}(a_{ik}^{\tilde{q}} P_R + b_{ik}^{\tilde{q}} P_L)\tilde{\chi}_k^0\tilde{q}_i + g\tilde{\chi}_k^0(a_{ik}^{\tilde{q}} P_L + b_{ik}^{\tilde{q}} P_R)\bar{q}\tilde{q}_i^*
\end{aligned}$$

- Quark-squark-gluino

$$\begin{aligned}
L_{q\tilde{q}\tilde{g}} &= -\sqrt{2}g_s T_{rs}^a \left[(\bar{q}_r P_R \tilde{g}^a \tilde{q}_{Ls} - \bar{q}_r P_R \tilde{g}^a \tilde{q}_{Rs}) + (\bar{g}^a P_L q_r \tilde{q}_{Ls}^* - \bar{g}^a P_L q_r \tilde{q}_{Rs}^*) \right] \\
&= -\sqrt{2}g_s T_{rs}^a \left[\bar{q}_r (R_{i1}^{\tilde{q}} P_R - R_{i2}^{\tilde{q}} P_L) \tilde{g}^a \tilde{q}_{is} + \bar{g}^a (R_{i1}^{\tilde{q}} P_L - R_{i2}^{\tilde{q}} P_R) q_r \tilde{q}_{is}^* \right]
\end{aligned}$$

- Gluon-gluino-gluino

$$L_{g\tilde{g}\tilde{g}} = \frac{ig_s}{2} f_{abc} G_\mu^a \bar{g}^b \gamma^\mu \tilde{g}^c$$

Squark-squark-gauge boson-gauge boson

$$\begin{aligned}
L_{\tilde{q}\tilde{q}\gamma\gamma} &= e^2 e_q^2 A_\mu A^\mu (\tilde{q}_L^* \tilde{q}_L + \tilde{q}_R^* \tilde{q}_R) = e^2 e_q^2 \delta_{ij} A_\mu A^\mu \tilde{q}_j^* \tilde{q}_i \\
L_{\tilde{q}\tilde{q}ZZ} &= \frac{g^2}{\cos^2 \theta_W} Z_\mu Z^\mu (C_{qL}^2 \tilde{q}_L^* \tilde{q}_L + C_{qR}^2 \tilde{q}_R^* \tilde{q}_R) \\
&= \frac{g^2}{\cos^2 \theta_W} Z_\mu Z^\mu (C_{qL}^2 R_{i1}^{\tilde{q}} R_{j1}^{\tilde{q}} + C_{qR}^2 R_{i2}^{\tilde{q}} R_{j2}^{\tilde{q}}) \tilde{q}_j^* \tilde{q}_i \\
&= \frac{g^2}{\cos^2 \theta_W} z_{ij} Z_\mu Z^\mu \tilde{q}_j^* \tilde{q}_i
\end{aligned}$$

Four squark interactions

$$\begin{aligned}
L_{\tilde{q}\tilde{q}\tilde{q}} &= -\frac{1}{2} g_s^2 T_{mn}^a T_{rs}^a \left(R_{i1}^\alpha R_{j1}^\alpha - R_{i2}^\alpha R_{j2}^\alpha \right) \tilde{q}_{jm}^{\alpha*} \tilde{q}_{in}^\alpha \left(R_{k1}^\beta R_{l1}^\beta - R_{k2}^\beta R_{l2}^\beta \right) \tilde{q}_{kr}^{\beta*} \tilde{q}_{ls}^\beta = \\
&= -\frac{1}{2} g_s^2 T_{mn}^a T_{rs}^a S_{ij}^\alpha S_{kl}^\beta \tilde{q}_{jm}^{\alpha*} \tilde{q}_{in}^\alpha \tilde{q}_{kr}^{\beta*} \tilde{q}_{ls}^\beta
\end{aligned}$$

KEY ADDITION FOR CARRYING SQUARK QCD WITH COMPLEX PARAMETER

The mix of boxers's arm moves and squark Then, the partial width of decay \tilde{q}_i ($\tilde{q}_i = \tilde{t}_i, \tilde{b}_i$) into final state fermions are

$$\begin{aligned}
\Gamma(\tilde{q}_i \rightarrow q + \tilde{\chi}_k^0) &= \frac{g^2 \lambda^{\frac{1}{2}} \left(m_{\tilde{q}_i}^2, m_q^2, m_{\tilde{\chi}_k^0}^2 \right)}{16\pi m_{\tilde{q}_i}^3} \times \\
&\left[\left(|a_{ik}^{\tilde{q}}|^2 + |b_{ik}^{\tilde{q}}|^2 \right) \left(m_{\tilde{q}_i}^2 - m_q^2 - m_{\tilde{\chi}_k^0}^2 \right) - 4 \operatorname{Re} \left(a_{ik}^{\tilde{q}*} b_{ik}^{\tilde{q}} \right) m_q m_{\tilde{\chi}_k^0} \right]
\end{aligned}$$

and

$$\begin{aligned}
\Gamma(\tilde{q}_i \rightarrow q + \tilde{\chi}_k^0) &= \frac{g^2 \lambda^{\frac{1}{2}} \left(m_{\tilde{q}_i}^2, m_q^2, m_{\tilde{\chi}_k^0}^2 \right)}{16\pi m_{\tilde{q}_i}^3} \times \\
&\left[\left(|a_{ik}^{\tilde{q}}|^2 + |b_{ik}^{\tilde{q}}|^2 \right) \left(m_{\tilde{q}_i}^2 - m_q^2 - m_{\tilde{\chi}_k^0}^2 \right) - 4 \operatorname{Re} \left(a_{ik}^{\tilde{q}*} b_{ik}^{\tilde{q}} \right) m_q m_{\tilde{\chi}_k^0} \right]
\end{aligned}$$

Partial width of decay \tilde{q}_i ($\tilde{q}_i = \tilde{t}_i, \tilde{b}_i$) of the final state boson (gauge and Higgs) would be:

$$\Gamma(\tilde{q}_i \rightarrow W^\pm + q'_k) = \frac{g^2 |A_{\tilde{q}_i \tilde{q}_j}^W|^2 \lambda^{\frac{3}{2}}(m_{\tilde{q}_i}^2, m_W^2, m_{q'_j}^2)}{16\pi m_W^2 m_{\tilde{q}_i}^3}$$

$$\Gamma(\tilde{q}_i \rightarrow Z + \tilde{q}_1) = \frac{g^2 |B_{21}^Z|^2 \lambda^{\frac{3}{2}}(m_{\tilde{q}_2}^2, m_Z^2, m_{\tilde{q}_1}^2)}{16\pi m_Z^2 m_{\tilde{q}_2}^3}$$

$$\Gamma(\tilde{q}_i \rightarrow H^\pm + \tilde{q}'_j) = \frac{g^2 |C_{\tilde{q}'_j \tilde{q}_i}^H|^2 \lambda^{\frac{3}{2}}(m_{\tilde{q}_i}^2, m_{H^\pm}^2, m_{q'_j}^2)}{16\pi m_{\tilde{q}_i}^3}$$

$$\Gamma(\tilde{q}_i \rightarrow H_i + \tilde{q}_1) = \frac{g^2 |C(\tilde{q}_1^\dagger H_i \tilde{q}_2)|^2 \lambda^{\frac{1}{2}}(m_{\tilde{q}_2}^2, m_{H_i}^2, m_{\tilde{q}_1}^2)}{16\pi m_{\tilde{q}_2}^3}$$

We have the following comments received on the results above. The process $e^+ e^- \rightarrow \tilde{q}_i \tilde{q}_j^*$ occurs via s-channel transmission with a photon and Z – boson particle

12. Applicability in practice: The estimated number may be partially verifiable credibility of the results obtained using experimental results from LEP, LEP2.

13. The following research: Results of calculation for all possible schema can conclude about the complexity of the parameters in the MSSM

14. All works published related essays: No

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Signature:

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